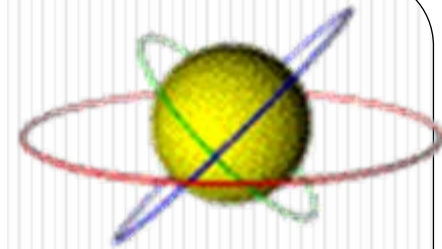




北京師範大學
BEIJING NORMAL UNIVERSITY



Orientation effects on evaporation residue cross sections in ^{48}Ca -induced hot fusion reactions

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2014.10



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2. 理论模型（双核系统模型）介绍

3. 对 ^{48}Ca 引起的热熔合的研究

3.1 熔合反应 $^{48}\text{Ca}+^{238}\text{U}$ 的蒸发剩余截面的方向效应研究

3.2 ^{48}Ca 引起的热熔合反应产生 $Z = 113$ 118的同位素的蒸发剩余截面计算

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4. 总结与展望



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冷熔合:

基于 ^{208}Pb 和 ^{209}Bi 靶

S. Hofmann *et al.*, Eur. Phys. J. A **14**, 147 (2002).

Z=111,112 **GAN**

K. Morita *et al.*, J. Phys. Soc. Jpn. **81**, 103201 (2012).

Z=113 **RIKEN**

热熔合:

Dubna

^{48}Ca 引起

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **70**, 064609 (2004)

Z=112

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **76**, 011601(R) (2007)

Z=113

Yu. Ts. Oganessian *et al.*, *Nature* **400** (1999)242; *PRC* **62** (2000)041604(R) ;
69 (2004)054607; **70** (2004)064609; *NPA* **734**(2004)109

Z=114

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **69**, 021601(R) (2004)

Z=115

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **63**, 011301(R)

Z=116

Yu. Ts. Oganessian *et al.*, Phys. Rev. Lett. **104**, 142502 (2010)

Z=117

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **74**, 044602 (2006)

Z=118

Yu. Ts. Oganessian *et al.*, Phys. Rev. C **79**, 024603 (2009).

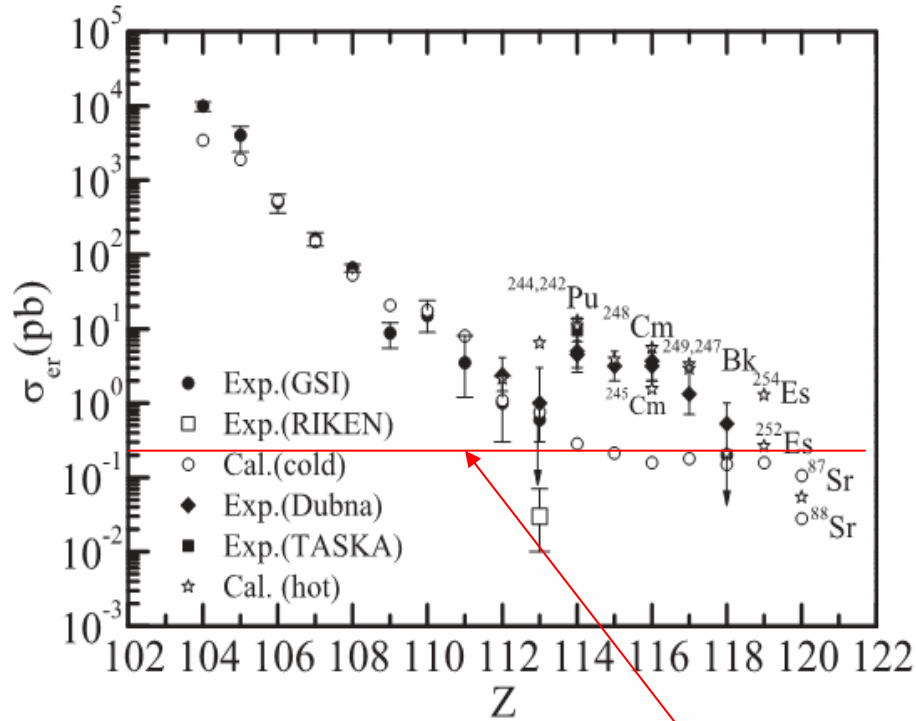
通过反应 $^{58}\text{Fe}+^{244}\text{Pu}$ 尝试合成**120**号元素

No decay chains consistent with fusion-evaporation reaction products were Observed. A more mass-asymmetry reaction would be preferable.

热熔和

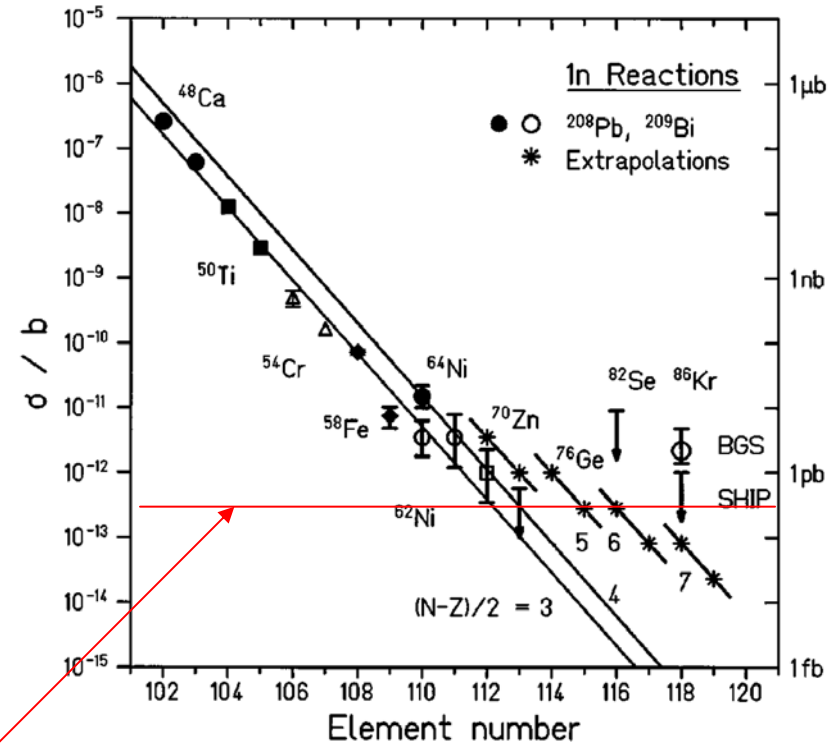
M. H. Huang *et al.*, *Phys. Rev. C* 82, 044614 (2010).

P. Armbruster, *Ann. Rev. Nucl. Part. Sci.*, 2000, **50**: 411



冷融合

S. Hofmann *et al.*, *Rev. Mod. Phys.*, 2000, **72**: 733



现在实验能探测的极限

1. 提高实验设备的探测灵敏度。
2. 找到更适合的反应体系或方法。





理论进展:

有很多模型和方法来研究超重核的合成机制:

DNS model:

G. G. Adamian, N. V. Antonenko, W. Scheid, and V. V. Volkov, Nucl. Phys. A 627, 361 (1997).

Zhao-qing Feng, Gen-ming Jin, Jun-qing Li, Nucl. Phys. A 836, 82 (2010).

Coupled Langevin-type equations:

V. Zagrebaev and W. Greiner J. Phys. G: Nucl. Part. Phys. 31(2005) 825–844

Fusion by diffusion model (FBD model): Smoluchowski diffusion equation

K. Siwek-Wilczynska, T. Cap, M. Kowal, *et al.*, Phys. Rev. C 87, 034616 (2013).

Zu-Hua Liu and Jing-Dong Bao, Phys. Rev. C 87, 034616 (2013).

Phenomenological approaches:

Ning Wang, Junlong Tian, and Werner Scheid, Phys. Rev. C 84, 061601(R) (2011).

L. Zhu, W. J. Xie, and F. S. Zhang, Phys. Rev. C 89, 024615 (2014).

R. Smolanczuk, Phys. Rev. C 81, 067602 (2010).

Nucleon collectivization mechanism:

V. I. Zagrebaev, Phys. Rev. C 64, 034606 (2001).

Two step model:

C. W. Shen, Y. Abe, D. Boilley, G. Kosenko, and E. G. Zhao, Int. J. Mod. Phys. E 17, 66 (2008).



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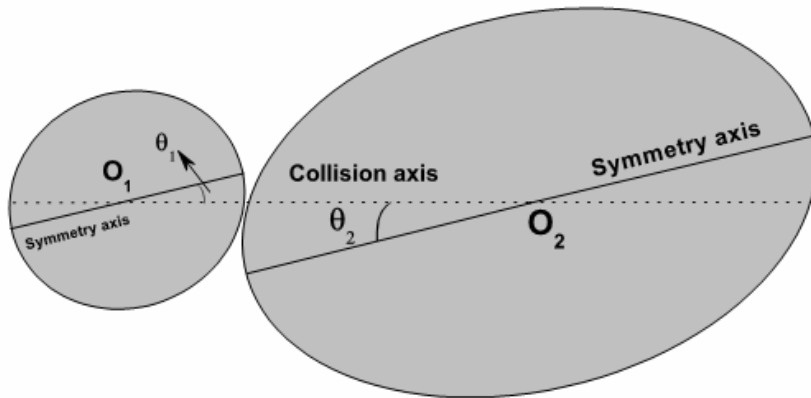


ER cross section:

$$\sigma_{\text{ER}}(E_{\text{c.m.}}) = \frac{\pi \hbar^2}{2\mu E_{\text{c.m.}}} \int_0^{\pi/2} \sin\theta_1 d\theta_1 \int_0^{\pi/2} \sum_J (2J + 1) T(E_{\text{c.m.}}, J, \theta_1, \theta_2) \times P_{\text{CN}}(E_{\text{c.m.}}, J, \theta_1, \theta_2) W_{\text{sur}}(E_{\text{c.m.}}, J) \sin\theta_2 d\theta_2$$

Capture cross section:

$$\sigma_{\text{cap}}(E_{\text{c.m.}}) = \frac{\pi \hbar^2}{2\mu E_{\text{c.m.}}} \sum_J (2J + 1) \times \int_0^{\pi/2} \int_0^{\pi/2} T(E_{\text{c.m.}}, \theta_1, \theta_2, J) \sin\theta_1 \sin\theta_2 d\theta_1 d\theta_2$$





俘获过程:

The transmission probability can be written as

$$T(E_{\text{c.m.}}, \theta_1, \theta_2, J) = \frac{1}{1 + \exp\left\{-\frac{2\pi}{\hbar\omega(\theta_1, \theta_2, J)}\left[E_{\text{c.m.}} - B(\theta_1, \theta_2) - \frac{\hbar^2}{2\mu R_B^2(\theta_1, \theta_2, J)}\right]\right\}}$$

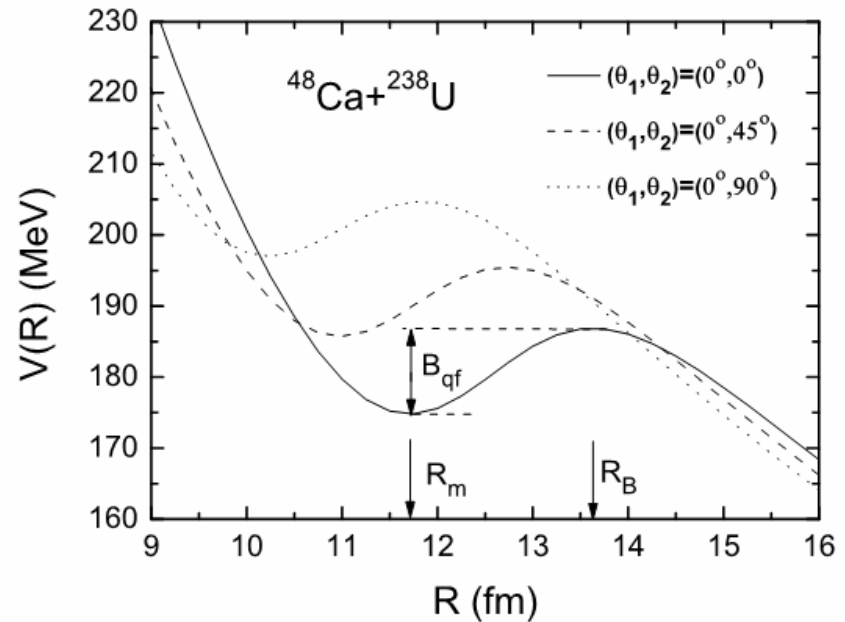
The nucleus-nucleus interaction potential

$$V(r, \theta_1, \theta_2) = V_N(r, \theta_1, \theta_2) + V_C(r, \theta_1, \theta_2)$$

$$V_C(r, \theta_1, \theta_2) = \frac{Z_1 Z_2 e^2}{r} + \left(\frac{9}{20\pi}\right)^{1/2} \left(\frac{Z_1 Z_2 e^2}{r^3}\right) \sum_{i=1}^2 R_i^2 \beta_2^{(i)} P_2(\cos\theta_i) + \left(\frac{3}{7\pi}\right) \left(\frac{Z_1 Z_2 e^2}{r^3}\right) \sum_{i=1}^2 R_i^2 [\beta_2^{(i)} P_2(\cos\theta_i)]^2.$$

$$V_N(r, \theta_1, \theta_2) = -V_0 \times$$

$$\left\{1 + \exp\left[\left(r - \sum_{i=1}^2 R_i [1 + (5/4\pi)^{1/2} \beta_2^{(i)} P_2(\cos\theta_i)]\right) a^{-1}\right]\right\}^{-1}$$





熔合概率: (DNS 模型)

$$P_{CN}(E_{\text{c.m.}}, J, \theta_1, \theta_2) = \sum_{A_1=1}^{A_{BG}} P(A_1, E_{\text{c.m.}}, J, \theta_1, \theta_2)$$

$$\frac{dP(Z_1, N_1, E_1, \theta_1, \theta_2, t)}{dt}$$

$$= \sum_{Z'_1} W_{Z_1, N_1; Z'_1, N'_1}(t) [d_{Z_1, N_1} P(Z'_1, N_1, E'_1, \theta_1, \theta_2, t)$$

$$- d_{Z'_1, N'_1} P(Z_1, N_1, E_1, \theta_1, \theta_2, t)]$$

$$+ \sum_{N'_1} W_{Z_1, N_1; Z_1, N'_1}(t) [d_{Z_1, N_1} P(Z_1, N'_1, E'_1, \theta_1, \theta_2,$$

$$- d_{Z_1, N'_1} P(Z_1, N_1, E_1, \theta_1, \theta_2, t)]$$

$$- [\Lambda_{\text{qf}}(\Theta(t), \theta_1, \theta_2) + \Lambda_{\text{fis}}(\Theta(t))] P(Z_1, N_1, E_1, \theta_1, \theta_2, t)$$

$$\Lambda_{\text{qf}}(\Theta(t), \theta_1, \theta_2) = \frac{\omega}{2\pi\omega^{B_{\text{qf}}}} \left[\sqrt{\left(\frac{\Gamma}{2\hbar}\right)^2 + (\omega^{B_{\text{qf}}})^2} - \frac{\Gamma}{2\hbar} \right] \\ \times \exp\left[-\frac{B_{\text{qf}}(Z_1, N_1, \theta_1, \theta_2)}{\Theta(t)}\right].$$

存活几率:

$$W_{\text{sur}}(E_{\text{CN}}^*, x, J) = P(E_{\text{CN}}^*, x, J) \times$$

$$\prod_i^x \left[\frac{\Gamma_n(E_i^*, J)}{\Gamma_n(E_i^*, J) + \Gamma_f(E_i^*, J)} \right]$$

G. G. Adamian, N. V. Antonenko, W. Scheid, and V. V. Volkov, Nucl. Phys. A 627, 361 (1997).

Zhao-qing Feng, Gen-ming Jin, Jun-qing Li, Nucl. Phys. A 836, 82 (2010).



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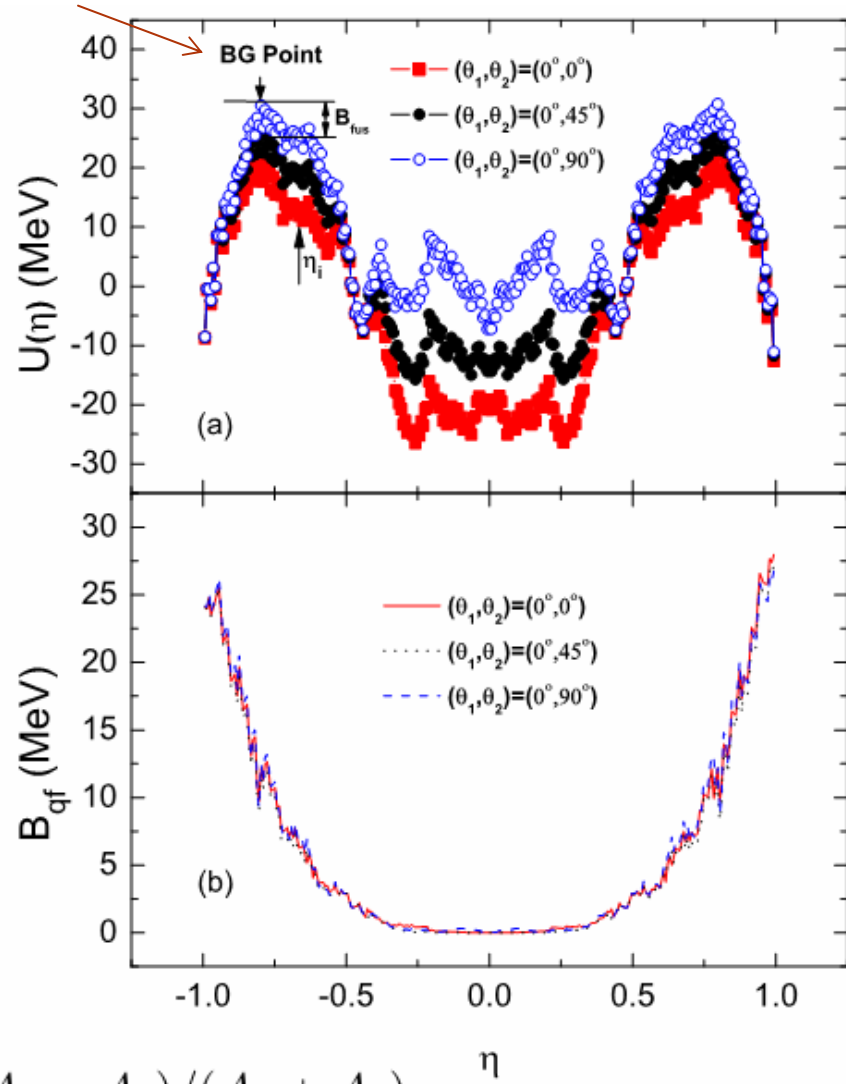
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Businaro-Gallone point $^{48}\text{Ca}+^{238}\text{U}$

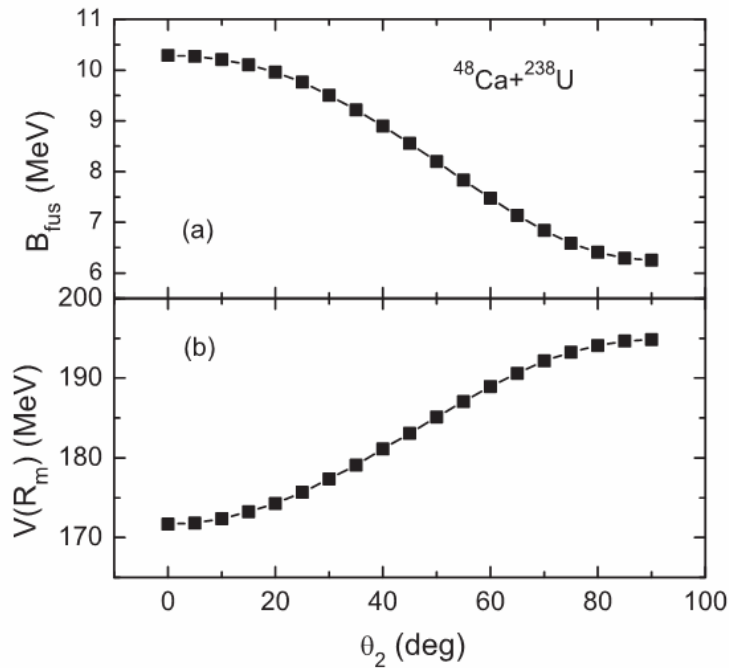
驱动势:

$$U(A_1, A_2, R, \beta_1, \beta_2, \theta_1, \theta_2, J) = U_{\text{LD}}(A_1) + U_{\text{LD}}(A_2) - U_{\text{LD}}(A) + V_{\text{CN}}(A_1, A_2, R, \beta_1, \beta_2, \theta_1, \theta_2, J)$$

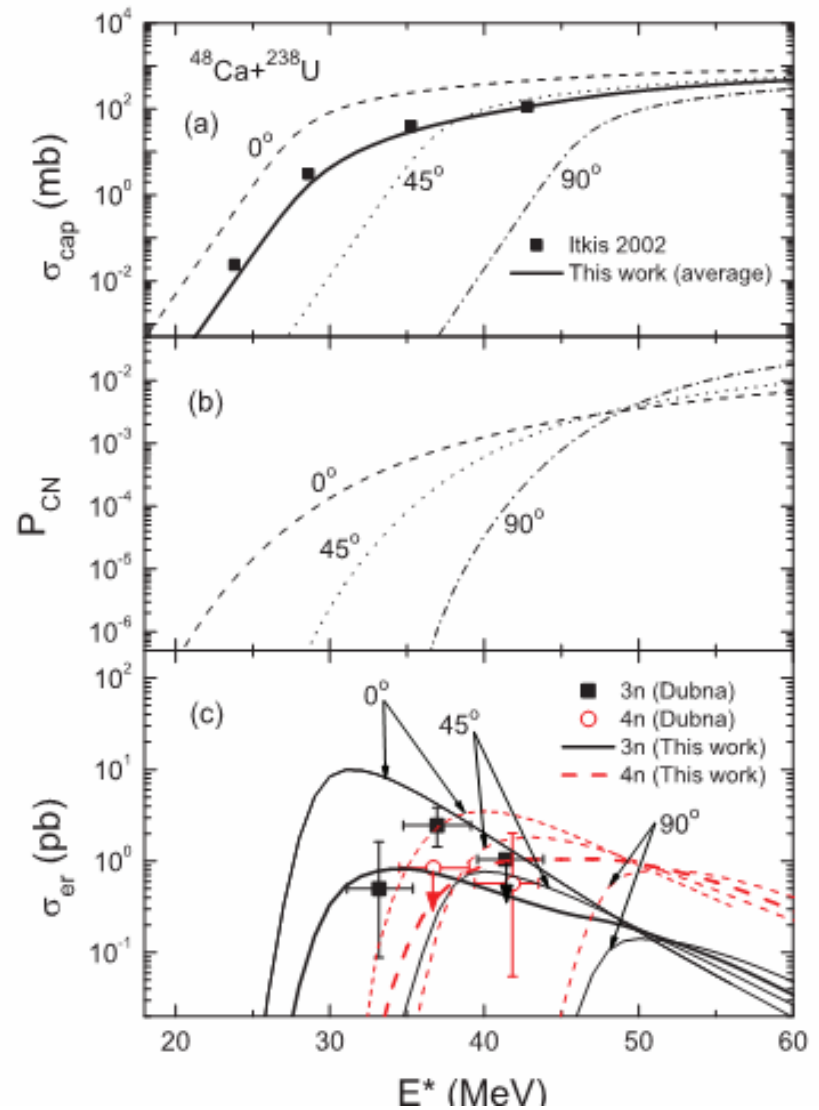


$$\eta = (A_1 - A_2)/(A_1 + A_2)$$

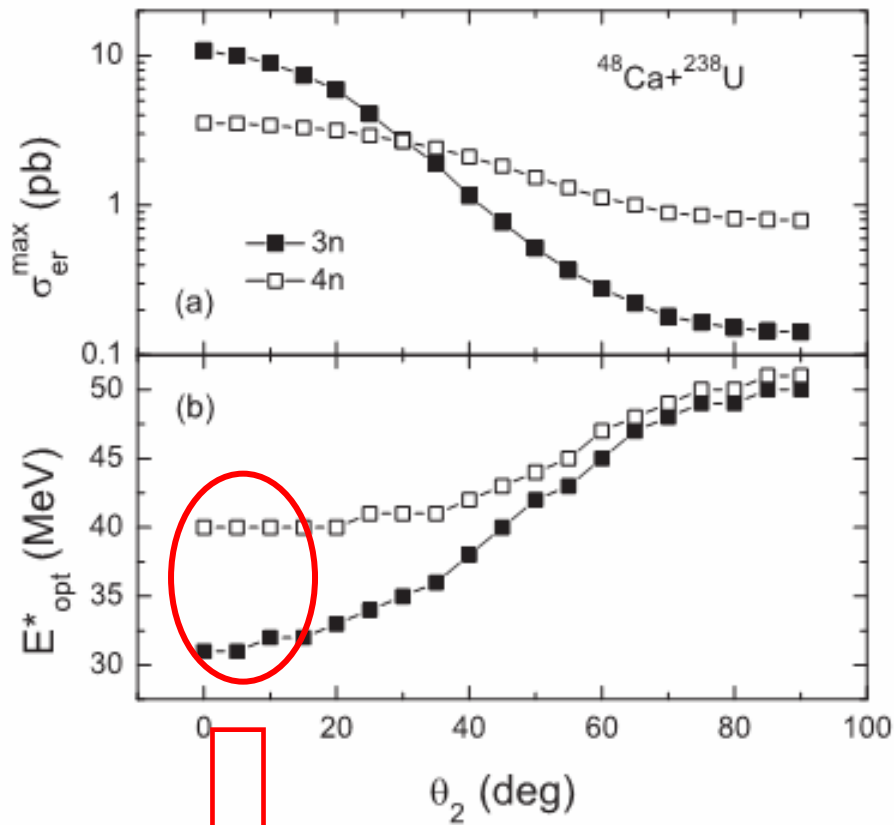
熔合反应 $^{48}\text{Ca}+^{238}\text{U}$ 的蒸发剩余截面的方向效应研究



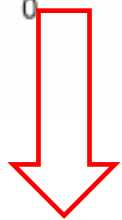
熔合概率的方向效应是内部熔合位垒 B_{fus} 和接触时的相互作用势 $V(R_m)$ 相互竞争的结果。



Study of $^{48}\text{Ca} + ^{238}\text{U}$ induced hot fusion reactions



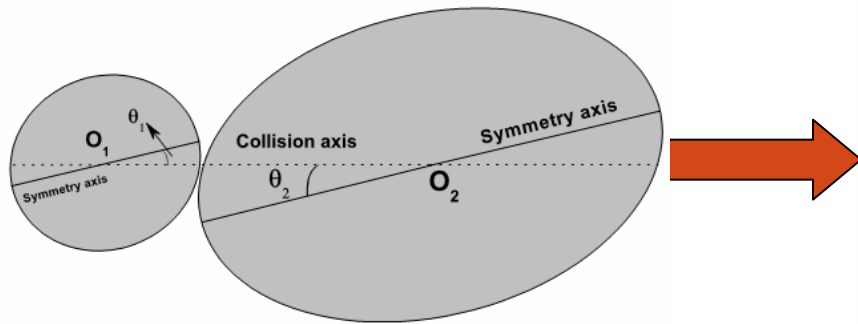
Decrease with the increasing value of theta2



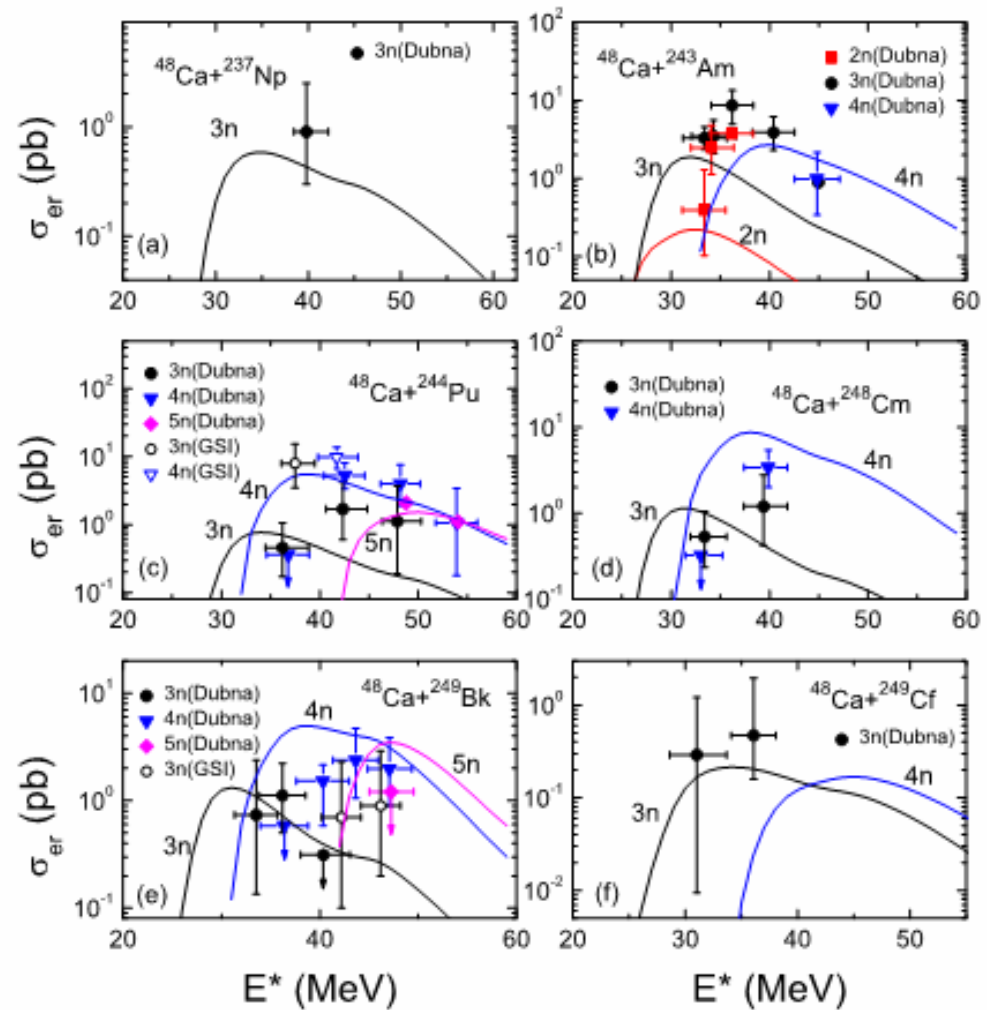
The discrepancy is large.

The maximal ER cross sections decrease with the increasing theta2 for both the 3n and 4n channel. The optimal excitation energies are increase by increasing theta2.

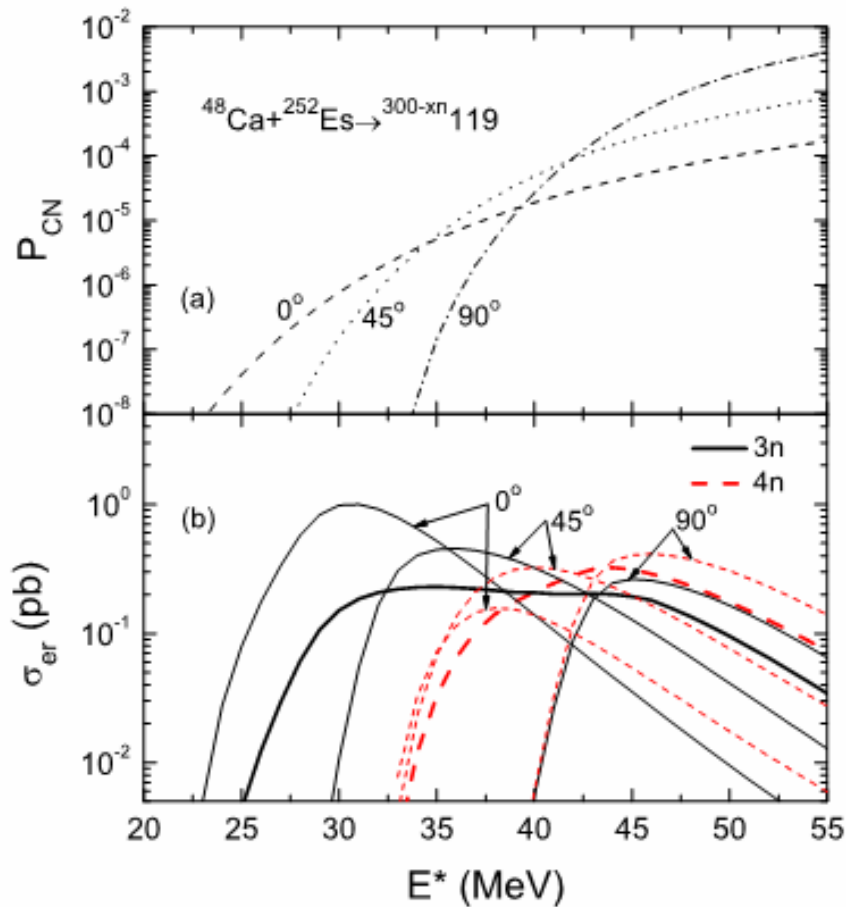
^{48}Ca 引起的热熔合反应产生 $Z = 113$ - 118 的同位素的蒸发剩余截面计算



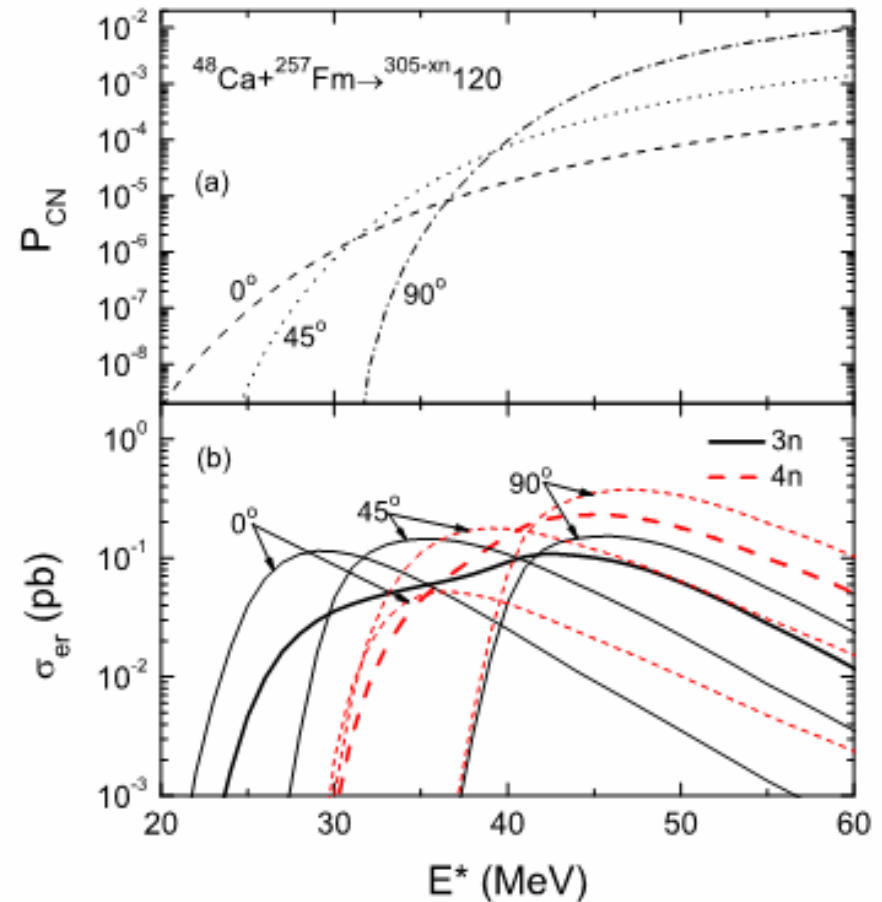
不同方向的平均结果



对超重核Z=119和120的预言



0.32 pb in the 4n channel



0.23 pb in the 4n channel

In the future, if it will be possible to prepare targets of ^{252}Es and ^{257}Fm , then the superheavy elements Z=119 and 120 probably can be synthesized.



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Conclusions

- The orientation of the colliding nuclei plays a significant role in fusion reactions.
- The capture cross sections for $\theta = 90^\circ$ are obviously suppressed, especially at sub-barrier energies. According to the DNS concept, the fusion probability strongly depend on the inner fusion barrier, the QF barrier, and the excitation energy of the DNS.
- It is found that fusion probability for $\theta = 0^\circ$ are much higher than those for $\theta = 90^\circ$ at lower incident energies, while the opposite behavior can be seen in the high incident region.
- The production cross sections of $Z = 119$ and 120 through the reactions $^{48}\text{Ca} + ^{252}\text{Es}$ and $^{48}\text{Ca} + ^{257}\text{Fm}$ are predicted. The maximal production cross sections of $Z = 119$ and 120 are 0.32 and 0.23 pb, respectively, in the $4n$ emission channel.

Thanks !